the quantum yield for this step is very high as compared with that of the first step (2).

The absorption peak of 1r at ca. 356 nm and that of 2 at ca. 300 nm are due to the charge transfer from the bridging peroxide to cobalt(III).<sup>23</sup> Such a charge-transfer excited state should be involved in the photodeoxygenation. The d-d excited state does not seem to lead to deoxygenation.

The quantum yield of **1r** on irradiation at 356 nm is relatively low as compared with those reported for the redox decomposition of some mononuclear cobalt(III) complexes upon irradiation of the ligand-to-metal charge-transfer band.<sup>24</sup> It is 1 or 2 orders of magnitude lower than those for the decomposition of ( $\mu$ superoxo)dicobalt(III) complexes.<sup>5,7</sup>

Despite the low quantum yield, the photoinduced deoxygenation is interesting, since it indicates the photocontrolled reversible oxygen uptake. The reversibility is unfortunately not maintained in the present systems due to the slow irreversible photodecomposition of the binuclear species to mononuclear cobalt(III) complexes.

2. Irreversible Decomposition of 1r to Mononuclear Cobalt(III) Complexes. 1r undergoes irreversible decomposition both thermally and photochemically, the thermal reaction being much slower and observed appreciably only at >50 °C. The mechanism of the thermal reaction is discussed first, since some relevant studies<sup>11,25,26</sup> permit more detailed discussion.

(a) The Thermal Reaction. A mononuclear (hydroperoxo)cobalt(III) complex was confirmed as an intermediate of the decomposition of  $[(CN)_5Co^{III}(\mu-O_2^{2-})Co^{III}(CN)_5]^{6-}$  in neutral solution<sup>25</sup> and of  $[(NO_2)(en)_2Co^{III}(\mu-O_2^{2-})Co^{III}(en)_2(NO_2)]^{2+}$  in acid solution.<sup>26</sup> Thus reaction 4<sup>27</sup> is likely to be the initial step of the decomposition of 1r.

$$[(en)_{2}Co^{III}(\mu-OH,O_{2}^{2-})Co^{III}(en)_{2}]^{3+} \rightarrow [Co^{III}(en)_{2}(H_{2}O)_{2}]^{3+} + [Co^{III}(en)_{2}(O_{2}H)(H_{2}O)]^{2+} (4)$$

In the absence of other ingredients, initially produced  $[Co^{III}(en)_2(O_2H)(H_2O)]^{2+}$  would slowly react with coordinated ligand

or ion-exchange resin during product analysis to give "orange species". Free ethylenediamine added to the solution of 1r would be oxidized by  $[Co^{III}(en)_2(O_2H)(H_2O)]^{2+}$ . The decrease in decomposition rate in the presence of free ethylenediamine is not clearly accounted for.<sup>28</sup>

In the presence of both free ethylenediamine and  $Co^{11}$  ions, reaction  $5^{27}$  would take place.

$$[Co^{III}(en)_{2}(O_{2}H)(H_{2}O)]^{2+} + 2[Co^{III}(en)_{n}(H_{2}O)_{6-2n}]^{2+} \rightarrow [Co^{III}(en)_{2}(H_{2}O)_{2}]^{3+} + 2[Co^{III}(en)_{n}(H_{2}O)_{6-2n}]^{3+} (5)$$
  

$$n = 2 \text{ or } 3$$

(b) The Photochemical Decomposition. The reaction products contain various unidentified species. The reaction may be initiated by homolytic cleavage of the  $Co^{III}$ – $(O_2Co^{III})$  bond as is the thermal reaction. The reactive intermediate,  $Co^{III}$ – $O_2^{2-}(H)$ , would undergo further complicated reactions. It is also plausible that the photodecomposition may be initiated by the reaction of 1 with  $Co^{II}$  species, which is produced by the initial photoequilibration. Addition of free ethylenediamine considerably decreases the amount of  $Co^{II}$  species by shifting the equilibrium to the dimer side, and it would retard the photodecomposition.<sup>28</sup>

Finally, it should be pointed out that previous studies on ( $\mu$ -peroxo)dicobalt(III) complexes with amines might have been affected by possible photochemical reactions and might require reexamination.

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(28) One possible explanation for the effect of free ethylenediamine on the rate of decomposition is that the free ligand would replace the hydroxo bridge to give the  $\mu$ -ethylenediamine complex,  $[(en)_2Co^{III}(\mu-O_2^{2^-},en)-Co^{III}(en)_2]^{4^+}$ , which has been claimed to exist (Crawford, M.; Bedell, S. A.; Patel, R. I.; Young, L. W.; Nakon, R. Inorg. Chem. 1979, 18, 2075-2079). The observed slow decomposition rate would be understood if the  $\mu$ -ethylenediamine complex were more stable toward a decomposition similar to reaction 4. Also the production of [Co<sup>III</sup>-(en)<sub>3</sub>]<sup>3</sup> <sup>+</sup> is expected from the decomposition of the  $\mu$ -en complex. The failure to observe any photodecomposition of 1r in the presence of free ethylenediamine would be explained if the postulated  $\mu$ -en complex were photostable. We are not fully convinced, however, of the existence of the  $\mu$ -en complex under our experimental conditions. No appreciable change in the absorption spectrum of 1r was observed on addition of free ethylenediamine, and no CD spectrum appeared on addition of (R)propylenediamine in the dark (the  $\mu$ -R-pn complex, if formed, would show some CD spectrum).

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# Hypervalent Tellurium in One-Dimensional Extended Structures Containing Te<sub>5</sub><sup>n-</sup> Units

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The electronic structure of three one-dimensional chains containing square-planar  $Te_5^{n-}$  units is presented, with employment of both molecular and band calculations. Two chains containing only  $Te_5^{2-}$  in different conformations and one in which they are modulated by Sn are treated. The  $Rb_2Te_5$  or  $Cs_2Te_5$  and  $K_2SnTe_5$  structures are thus modeled. A three-center, four-electron model can be used to understand the elongation of the Te-Te bands within the  $Te_5$  unit, whose inherent instability is overcome by the formation of the chains. It is suggested how oxidation of the chain, for instance, by suitable choice of synthetic conditions, would alter the geometric physical properties of these polymers.

The recent literature contains a number of examples of a square-planar tellurium structural unit which may formally be defined as  $Te_5^{\prime\prime}$  (1). In some cases such as  $Rb_2Te_5^1$  and  $Cs_2Te_5^2$ 



the stoichiometry clearly defines the charge on the unit as 2-. In these instances the unit is the basic building block of a one-di-

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<sup>(27)</sup> The products of reactions 4 and 5 should be in equilibrium with their conjugate bases, hydroxo complexes, depending on the pH of the solution.

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mensional anionic chain. In Cs<sub>2</sub>Te<sub>5</sub> the Te<sub>5</sub>" units are screw axis related (2) while in  $Rb_2Te_5$  the  $Te_5^{n-1}$  units are related by translation (3).



For most of the other cases in which the  $Te_5^{n-}$  unit appears, the attribution of a particular charge to the unit is more ambiguous. For instance, on  $K_2SnTe_5$ ,<sup>3</sup> the unit again appears embedded in a one-dimensional chain (4) but the Te<sub>5</sub><sup>*n*</sup> units alternate



with tetrahedral Sn in making up the chain. Formally, at least, the square-planar units could still be considered  $Te_5^{2-}$  if the tin is Sn(0), but the Te-Sn bond length of 2.74 Å is almost identical with the sum of the tetrahedral covalent radii (2.72 Å) given by Pauling.<sup>4</sup> This suggests a formal oxidation state of 4+ for the Sn and a net charge of 6- on the  $Te_5^{n-}$  unit.

A similar ambiguity arises in the case of Ga<sub>2</sub>Te<sub>5</sub>,<sup>5</sup> although this is no longer a linear-chain structure but a three-dimensional one, as depicted in 5 in a view down the c axis of the tetragonal cell.



The Te<sub>5</sub> units are clearly present, with Te(1) as the central atom and four Te(2) atoms in the square-planar arrangement. The gallium is tetrahedrally surrounded by four Te(2) atoms. A choice of the "normal" valence state of Ga(III) leads to a formal charge of 6- on Te<sub>5</sub><sup>*n*</sup> while a choice of 2- for the charge on Te<sub>5</sub><sup>*n*</sup> would lead to an assignment of the only rarely encountered Ga(I).<sup>6</sup>

The structure of  $AgTe_3^7$  (6) may also be visualized as being built up from  $Te_5^{n-}$  units (one such unit has been darkened) to

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Table I. Geometric Features of the Te<sub>5</sub><sup>n-</sup> Unit

	<i>a</i> , Å	b, Å	c, deg	d, deg			
Cs <sub>2</sub> Te <sub>5</sub>	3.046 (1)		90.5	89.5			
Rb <sub>2</sub> Te <sub>5</sub>	3.039(1)		92.0	88.0			
$K_2 SnTe_5$	3.015 (5)	3.052 (5)	95.0	88.2, 88.8			
Ga,Te,	3.027 (2)		90.0				
AgTe	3.052 (5)		89.8	89.8, 90.5			
NdTe <sub>1</sub>	3.076 (10)		90				
$Re_2Te_5$	2.983 (3)	3.022 (3)	91.2	88.8			

give a rhombohedral structure, which is derived from the simple cubic  $\alpha$ -polonium structure by a small perturbation.



The structures of NdTe<sub>2</sub><sup>8</sup> and NdTe<sub>3</sub><sup>9</sup> also contain identifiable units of Te<sub>5</sub><sup>n-</sup> where assignment of a formal charge is virtually ruled out by the complexity of the structures. Part of the NdTe<sub>3</sub> structure is shown in 7.



Finally, the fascinating structure of Re<sub>2</sub>Te<sub>5</sub> has been reported by Klaiber, Petter, and Hulliger.<sup>10</sup> The structure is composed of [Re<sub>6</sub>Te<sub>8</sub>]<sup>2+</sup> clusters linked by butterfly-like units formally written as  $\{[]Te_6]Te$ , where the last-noted Te is the central one of the unit (8). The  $Te_5^{n-}$  unit is easily recognized here, however,



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Figure 1. Energy level diagram for ICl<sub>4</sub><sup>-</sup> at two different geometries.

and exhibits geometric features characteristic of other units we have already noted.

What are these features? They are summarized in Table I. The variation in geometry among these units is not large, considering the differences in crystalline environment in which they are found. The average Te-Te bond length is 3.039 Å, and except for the one large angle in K<sub>2</sub>SnTe<sub>5</sub> they are all essentially square planar in geometry. This characteristic bond length is significantly longer by 0.25–0.30 Å than twice the covalent radius<sup>4</sup> of 1.37 Å, since normal covalent Te-Te bonds fall in the range 2.69-2.80 Å.11

The essential questions we wish to address here, then, are as follows: how can the formation of the chain compounds be visualized, and why are the Te(central)-Te(peripheral) bond lengths longer than normal Te-Te bond lengths by  $\sim 0.25$  Å? We note that the same questions have been recently analyzed in a paper by Bullett.<sup>12</sup>

We turn to the second question first, by recalling an observation we made earlier regarding the  $\{SnTe_5^{2-1}\}_{\infty}$  anionic chain. Considering the building blocks of this polymeric chain as  $Te_5^{6-}$  and Sn(IV) may at first glance lack chemical aesthetic appeal. However, by analogy with some other well-known species such an approach can lead to an understanding of the bonding in this system.

The square-planar unit is reminiscent of mixed-halogen anions  $BrF_4^-$  and  $ICl_4^-$  and of  $XeF_4$ , an observation already made by others.<sup>2,10</sup> Let us initially examine the first two members of this group. Both are well-documented square-planar structures,<sup>13-15</sup> in which the Br-F bond length is 1.89 (2) Å while I-Cl ranges from 2.42 (1) to 2.60 (1) Å and averages 2.51 (2) Å.<sup>16</sup> These distances are again significantly longer than the respective single-bond distances of Br-F (1.756 Å) and I-Cl (2.321 Å) in the diatomic molecules.<sup>17</sup> This is the same lengthening observed in the  $Te_5^{n-}$  units, and its origin is an important point of our story.

The bonding in these molecules is clearly best described in terms of the classical three-center-four-electron bonding model.<sup>18,19</sup> We

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Figure 2. Walsh diagram for the  $D_{4h}-C_{2v}$  transformation of the Te<sub>5</sub> unit. will now discuss in detail the electronic structure of square-planar

 $ICl_4$  and  $Te_5^{n-1}$ . The energy level diagram for  $ICl_4^-$  in  $D_{4h}$  symmetry is given in Figure 1. The two highest occupied MO's in this 36-electron system are both antibonding  $(a_{1g} \text{ and } a_{2u})$  and are clearly destabilizing at a normal I-Cl bond length of 2.32 Å. Their influence, especially that of the  $a_{1g}$ , is significantly decreased by extending the bond length to 2.505 Å. Formation of the anion from I<sup>-</sup> and four Cl atoms leads to a net decrease in energy of 48.8 kcal/mol per I-Cl bond formed, a value that compares favorably with the literature value of 49.7 kcal/mol.<sup>20</sup> It would appear that lengthening of the I-Cl bonds is required to reduce the antibonding character of the HOMO. An analogy could, in principle, be made with SF<sub>4</sub>, which has the  $C_{2v}$  trigonal-bipyramid structure with one lone pair of electrons in an equatorial orbital; it is a 34-electron system. The Walsh diagram for the  $D_{4h} \rightarrow C_{2v}$  transformation for a Te<sub>5</sub> unit with Te-Te bond lengths of 2.91 Å is given in Figure 2. For a 32-electron count for  $D_{4h}$  Te<sub>5</sub><sup>2-</sup> the e<sub>g</sub> level is the HOMO, while for the 36-electron count the  $a_{1g}$  level is the HOMO. The filled  $a_{2u}$  and  $a_{1g}$  orbitals of Figure 1 are not shown in Figure 2. They are just below the  $e_g$  level. In the 32-electron case the  $C_{2v}$ geometry is more stable by about 0.45 eV. Filling the  $a_{2u}$  level for a 34-electron count raises the total energy of the  $C_{2v}$  structure by 0.54 eV, barely favoring the  $D_{4h}$  geometry. For the 34-electron systems, e.g.  $SF_4$ , there is a delicate balance in the calculations between planar and  $C_{2v}$  structures.<sup>19</sup> However, the 36-electron count raises the energy of the  $C_{2v}$  above the  $D_{4h}$  by an additional 1.43 eV, so the latter geometry is clearly preferred.

Let us now carry the analogy directly over to the  $Te_5^{n-}$  unit. First, we must take care of the electron count, which again will be 36 if n = 6. The formal oxidation state on the central Te is II, considering each of the peripheral Te's as 2-.

When the Te-Te distance is increased from 2.84 to 3.04 Å, the  $a_{1g}$  level drops by 2.95 eV and the  $a_{2u}$  level drops by 0.4 eV while other levels remain essentially constant. Thus the antibonding character is significantly decreased with almost no loss in bonding character. So why does the Te<sub>5</sub><sup>6-</sup> unit not exist as an independent entity like ICl<sub>4</sub><sup>-</sup> and its analogues? The reason is the difference in electronegativity in the two atom types for  $ICl_4^-$  and the lack

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whereas, as expected, there is much less polarization in the  $Te_5^{6-}$  case. Thus, in spite of the bond lengthening there is still sufficient antibonding character between central and peripheral Te's to destabilize an individual unit.

There is another factor operating to make  $XY_4$  hypervalent molecules kinetically unstable. Orbitals such as 9 or 10 have substantial electron density on the external Y ligands and are X-Y antibonding. The less electronegative Y is relative to X, the more will such orbitals be localized on the external atoms; these in turn will be activated with respect to electrophiles. The general prescription for kinetic stability for XY<sub>4</sub> molecules of this type is to make the external Y groups as electronegative as possible. In lieu of a large difference in electronegativity electron "pileup" on terminal atoms can be alleviated by forming interfragment bonds.

For Te<sub>5</sub><sup>6-</sup> stabilization can be achieved by the formation of bonds between units, as in the chain compounds 2 and 3, or between Te<sub>5</sub><sup>6-</sup> units and a suitable "mediator" such as Sn. Let us demonstrate how this happens, by first building model molecules from these Te<sub>5</sub><sup>6-</sup> units, and then extending them to the infinite chains observed.

Let us make the bonding in  $Te_5^{6-}$  explicit by structure 11, which shows the electron pairs as bars. The central Te has four electrons



in this extreme resonance structure and uses two formally empty p orbitals to engage in electron-rich three-center bonding with the external tellurium atoms. Please do not consider the crude assignment of electrons in **11** as anything but an extreme formalism.

Where are the frontier orbitals 9 and 10,  $a_{1g}$  and  $a_{21}$ , in this picture? They are to be identified with the two lone pairs at the central atom in 11, as approximate as these identifications perforce must be.

The rough formalism illustrated in structure 11 does indicate how a  $Te_5^{6-}$  fragment can bond to other entities. If it retains all of its electrons, it can only act effectively as a Lewis base, through either the terminal or central atom lone pairs. This is what it does in 4, the formal Sn(IV) complex of  $Te_5^{6-}$ .

If we allow oxidation of the  $Te_5^{6-}$  unit, further bonds may be formed. For instance, one view of the  $Te_5^{2-}$  system obtained by removal of four electrons from  $Te_5^{6-}$  is shown in 12. We have





Figure 3. Interaction diagram for 13 and 14. Only those orbitals on the fragments that interact have been drawn and labeled. The HOMO's for the fragments are also indicated.

chosen to oxidize the terminal tellurium atoms. 12 is a tetraradical, set up to polymerize with itself to form Te-Te bonds and extended structures of type 2 and 3.

We now return to another approach to the electronic structure of the polymeric  $(Te_5^{2-})_n$  chains by examining the bonding in two model compounds, 13 and 14. We have termed them cisoid and



transoid, in view of their topological similarity to *cis*- and *trans*-polyacetylene. In the  $Rb_2Te_5$  structure Te-Te bonds within a square-planar unit are 3.04 Å, while those between units are 2.76 Å. An intermediate value of 2.91 Å was chosen for all. All bond angles were chosen as 90° around a central Te atom, while 102.5° was used for the angle around the peripheral atoms.

The entire unit is assigned a charge of 2–. One way to reach this 72-electron count is to begin with a 96-electron trimer  $(Te_5^{2-})_3$ , which is terminated by four H<sup>-</sup> groups instead of Te<sup>2–</sup>. The 24 lost electrons are essentially the lone pairs on the replaced Te atoms. In an analysis of the bonding in 13 or 14 let us first consider a fragmentation scheme wherein the central Te<sub>5</sub> unit is one fragment and the two peripheral units comprise the second fragment. Again, all Te–Te bond lengths were given an intermediate value of 2.91 Å. Te–H was taken as 1.72 Å.

The interaction diagram is shown in Figure 3. The interaction of the two terminal  $Te_3H_2$  groups with the central  $Te_5$  results in the formation of four Te-Te  $\sigma$  bonds, from the interaction of essentially  $p_z$  orbitals. The antibonding combinations are pushed above the  $a_1$  orbital derived from the  $a_{1g}$  level of the Te<sub>5</sub> unit, which now becomes the HOMO for the system. This and the second highest filled orbital, also  $a_1$ , are both antibonding within the Te<sub>5</sub> unit. The interfragment bonding combinations are not stabilized to the same extent as the destabilization of the intrafragment



Figure 4. Interaction diagram for 16.

antibonding combinations. The HOMO's for the fragments are at approximately -10.6 eV while the HOMO for the full molecule is at -6.9 eV. As a result of this large difference, the net effect of forming the molecule is a destabilization of the structure with respect to the two fragments by 1.66 eV. The overall situation may be described as the formation of four  $\sigma$  bonds between units at the expense of filling antibonding orbitals within a unit.

The net positive energy for this process still leaves us with no rationalization for the formation of the compound. Recall that we chose an intermediate value of 2.91 Å for all the Te(peripheral)-Te(peripheral) bond lengths. If we now increase the Te-(peripheral)-Te(central) bond length to that observed in the crystal (3.034 Å) and likewise reduce the Te-Te bond length between square-planar units to 2.84 Å, we obtain a net stabilization of 2.2 eV, or about 12.6 kcal per Te-Te bond. The stabilization is due principally to a lowering of the HOMO (by 1.3 eV). This rather dramatic change with geometry demonstrates the strong competition between the antibonding character of the Te<sub>5</sub><sup>6-</sup> unit discussed earlier and the stabilization obtained by forming chains of these units.

The difference in the nature of the Te-Te bonds in this model for a polymer was already foreshadowed in the electronic structure of the initial model, with all Te-Te distances equal. The overlap populations computed there were 0.155 within the square-planar unit and 0.449 between units. This is a large difference, reflecting the weaker electron-rich three-center bonding. It is no wonder that the bonds to the square-planar coordinated Te elongate.

Further evidence for this driving force for Te(central)-Te-(peripheral) bond elongation may be derived from a calculation in which the fragmentation is chosen to examine the effect of bonding within these square-planar units, rather than between them, in the stabilization of the chain structure. One fragment would be the four ditelluride (Te-Te) units that make up the links between square-planar units, while the remainder of the molecule (two TeH<sub>2</sub> units plus Te) comprises the second fragment. For the geometry with all Te-Te bonds of 2.91 Å, the bringing together of these fragments (formation of bonds within square-planar Te<sub>5</sub> units) is unstable with respect to the fragments by 3.52 eV. Recall that relaxation of the geometry to that found in the crystal for the square-planar fragmentation did result in a net stabilization of 2.2 eV. The same process for this fragmentation still does not suggest a net bonding situation—the energy of the whole is 1.24 eV higher than the sum of the fragments.

In summary, the square-planar Te<sub>5</sub><sup>6-</sup> system can find a way of overcoming a delicately balanced but essentially antibonding situation by losing four electrons and by forming bonds between neighboring units. Now that we have seen the basis for the formation of bonds between square-planar units, we wish to ask if the geometry also plays an important role in this process, for we have already seen that even a homoatomic chain can adopt at least two different geometries, 13 and 14. For the transoid model compound, 14, the interaction diagram is virtually identical with that for the cisoid, 13, already discussed. This is not unreasonable: the structure is stabilized by the formation of  $\sigma$  bonds between essentially  $D_{4h}$  units, and the propagation of these bonds "above" or "below" the plane would seem to be equally possible on an energetic basis. The only real difference in interatomic interactions between the two molecules is a possible through-space Te---Te interaction indicated in 15. But the distance is nearly 6 Å, so this may be ignored.



We are now in a position to look at the  $SnTe_5^{2-}$  chain, 4, and again we choose a model compound, 16, to get some idea of the bonding in this system. We suggested earlier two possible ways



of thinking about the electron count: either as a combination of  $\mathrm{Sn^0}$  and  $\mathrm{Te_5}^{2-}$  units or as  $\mathrm{Sn(IV)}$  and  $\mathrm{Te_5^{6-}}$ . In light of what has been said above, clearly there is a preference for the second choice.

Calculations were carried out for two geometries, 17a and 17b. The restrictions of 90° bond angles around the central Te and Te-Sn bond lengths of 2.78 Å determine the Te-Te-Sn and Te-Sn-Te bond angles. In the experimental structure the Te-



Te-Te angles within a Te<sub>3</sub>Sn ring are 88.5 and 95.0° while exocyclic ones are 88.5°. The angles about the Sn are 102.1, 103.9, and 112.8° (4×) so strict square-planar geometry is not obtained in the former case and ideal tetrahedral geometry is not obtained in the latter case.

The interaction diagram for a fragmentation into planar Te<sub>5</sub> (and two planar Te<sub>3</sub>H<sub>2</sub> units) in one fragment and two Sn in the second fragment for geometry 17a is given in Figure 4. A pattern similar to that which we encountered earlier for 13 is observed here as well. Eight bonding orbitals are only slightly lowered with respect to their energy in the Te<sub>5</sub> and Te<sub>3</sub>H<sub>2</sub> fragments, but the antibonding orbitals are pushed way above the by now familiar antibonding a<sub>1</sub> orbital within the central Te<sub>5</sub> unit, which becomes the HOMO. The lowest antibonding combination can be found at +0.59 eV. Note that the HOMO is lowered by ~0.25 eV upon formation of the molecule due to a small contribution from the p<sub>x</sub> orbital of the Sn, which leads to some Te(peripheral)-Te(peripheral) bonding character. When the geometry is altered to 17b, increasing the bonded Te-Te distance and decreasing the inter-



Figure 5. Band structure for the transoid  $Te_5^{2-}$  chain.

action between peripheral Te's through the Sn in the HOMO, the HOMO is lowered by 1.25 eV, indicating again the importance of the antibonding character within this unit. However, the block of bonding orbitals at  $\sim 11.5$  eV does not contain any contribution to bonding within Te<sub>5</sub> units.

The formation of the molecule from these two fragments with geometry 17a yields a net binding energy of 15.3 eV or 44 kcal per Sn-Te bond formed. Relaxation to the geometry of 17b yields 18.5 eV or 53 kcal per Sn-Te bond. The molecule may also be fragmented as in 18 to examine the effect of the formation of



Te(central)-Te(peripheral) bonds. With a charge of 2- on the central Te the net binding energy (for the formation of eight Te-Te bonds) is 0.56 eV, for the geometry in 17a. Clearly the first fragmentation scheme with the geometry in 17b is preferable in accounting for the bonding in this model compound.

We are now ready to examine the extended systems containing Te<sub>5</sub> units. The band structure for the transoid Te<sub>5</sub><sup>2-</sup> chain 3 is shown in Figure 5. The Fermi level is at -6.95 eV (k = 0) corresponding to the energy of the HOMO in the trimer 14. There is very little dispersion in this band due to the near-orthogonality of the now familiar orbitals on the peripheral Te's at both edges of the Brillouin zone.

Two bonds are formed per unit cell (or more correctly between unit cells as defined here), so by analogy with the molecular unit we expect two essentially antibonding levels above the Fermi level; these appear at -5.71 and -5.33 eV, respectively. The material is thus expected to be a semiconductor with band gap 1.24 eV at k = 0.



Figure 6. Density of states for the transoid  $Te_5^{2-}$  chain.

Figure 5 shows only those bands that are important in Te-Te bonding. Omitted from the figure, for reasons of clarity, are several bands between -9.8 and -13.5 eV, which are composed mainly of nonbonding orbitals. The important interactions at the edges of the Brillouin zone are summarized in Figure 5. The antibonding orbitals appear in symmetric and antisymmetric combination with respect to the mirror plane of the chain as in 19.



Then, we expect two bonding bands, again corresponding to the A and S combinations in **20**. In fact, four bands may be associated with the bonding due to rather extensive mixing between bands at the edges of the Brillouin zone. The highest S antibonding band at k = 0 (7) correlates with the highest bonding band at k = 0.5 (4'). The latter is, in turn, antibonding with respect to its origin at k = 0 due to the presence of antibonding interactions within the Te<sub>5</sub> unit. The symmetric antibonding band at k = 0.5 (2).

As one starts from k = 0.5 for the S antibonding band (6') there is a correlation with the lower S bonding band at -12.1 eV (2). These correlations from antibonding to bonding bands lead to avoided crossings which account for the change in topology of the band orbitals in moving across, e.g., a bonding band.

A similar correlation may be made for the antisymmetric bands. The antibonding one at k = 0 (6) correlates with the higher bonding one at k = 0.5 (3') while the lower A band at k = 0 (1) correlates with A antibonding at k = 0.5 (7').

Associated with the band structure of the transoid chain is the density of states (DOS) in the polymer. This is plotted in Figure 6, along with a partitioning that shows the contribution of the four peripheral telluriums. The remainder in the DOS curve is the contribution of the single central Te. The total charge on a peripheral Te atom is -0.587, and that on the central Te is +0.348.



Figure 7. Crystal orbital overlap population (COOP) curves for the two different Te-Te bonds in the  $Te_5^{2-}$  chain.



Figure 8. Band structure of the cisoid  $Te_5^{2-}$  chain. Only bands contributing to  $Te_5$ ---Te<sub>5</sub> interaction are included. The two-letter symmetry classifications refer to the mirror plane parallel to the chain propagation direction and the  $2_1$  screw axis.

This is the same trend as in the Te<sub>5</sub> model.

The total overlap populations in the polymer are 0.155 for  $Te_c-Te_p$  and 0.429 for  $Te_p-Te_p$ , where the subscripts c and p stand for central and peripheral, respectively. A more detailed decomposition of the bonding is obtained through COOP curves. These are overlap-population-weighted densities of states, i.e. DOS curves in which the relative number of states in a given interval is weighted by the contribution that those states make to bonding or antibonding for a specified bond. We have found these COOP curves most useful for analyzing bonding trends.<sup>21</sup> A COOP curve for the  $Te_5^{2-}$  polymer is shown in Figure 7. Note that the highest occupied band in the polymer is both  $Te_p-Te_p$  and  $Te_c-Te_p$  antibonding, especially so in the latter bond. The antibonding nature of the model and the polymer has been discussed earlier. It is interesting to speculate here that a material with stronger Te-Te bonding could be made if one could oxidize the  $Te_5^{2-}$  chain. For



Figure 9. Band structure for the  $SnTe_5^{2-}$  chain. The two-letter symmetry classifications refer to the mirror plane parallel to the chain propagation direction and the  $4_2$  screw axis. E refers to a degenerate pair of bands.



Figure 10. Total density of states and contributions from the three different atom types in the  $SnTe_5{}^{2-}$  chain.

instance, the total overlap for the  $Te_c-Te_p$  and  $Te_p-T_p$  bond rises to 0.373 and 0.471, respectively, upon reducing the electron count from 32 to 30 electrons per Te<sub>5</sub> unit.

The similarity between the model cisoid and transoid molecules 13 and 14 is expected to be maintained in the extended structures as well. The major difference is the presence of the additional screw-axis element in the cisoid chain, causing the repeat distance to double. The symmetry of the chain is thus nonsymmorphic, resulting in a "folding back" of the bands, as seen in Figure 8. Again, only those bands that contribute to bonding between Te<sub>5</sub> units have been included. The topology of the orbitals is the same for both structures. Indeed, if we could "unfold" the bands for the cisoid polymer, the two band structures would be essentially superimposable.

The band structure for the  $SnTe_5^{2-}$  chain (Figure 9) shows even less dispersion than the pure Te<sub>5</sub> chains. The band gap in this case is 2.66 eV, again indicating a semiconductor. Following the model molecule 16, the conduction band (~-5.5 eV) is composed of nonbonding orbitals on the tellurium atoms. The band structure has been simplified here by eliminating all nonbonding orbitals except those in the conduction band. A discussion similar to that given for the pure Te<sub>5</sub> chains relating to the avoided crossings could

<sup>(21)</sup> These curves were invented by S. Wijeyesekera in our group. For some applications see: (a) Hughbanks, T.; Hoffmann, R. J. Am. Chem. Soc. 1983, 105, 3528-3537. (b) Kertesz, M.; Hoffmann, R. Ibid. 1984, 106, 3453-3460. (c) Saillard, J.-Y.; Hoffmann, R. Ibid. 1984, 106, 2006-2026.



Figure 11. COOP curves for the two bond types in the  $SnTe_5^{2-}$  chain.

be given here; however, the bonding in the chain can be understood on the basis of the model compound 16 discussed above.

The total density of states of  $SnTe_5^{2-}$  is shown in Figure 10. Note the relatively small contribution of Sn to the DOS; this is consistent with the large positive charge calculated on Sn of 1.13 and the formal assignment of a +IV oxidation state. The highest occupied band is almost entirely on Te. It represents mainly the nonbonding electrons that have been depicted earlier in Figure 7. There are some differences between our band structure and that computed by Bullett,<sup>12</sup> but there are more similarities: the absence of Sn character in the highest band and the concentration of Sn in a band ~7 eV below the Fermi level.

The COOP curves in Figure 11 reflect the pattern already discerned in the model compound: below the Fermi level the interaction between peripheral Te and Sn is all bonding; the Te(central)-Te(peripheral) overlap populations are bonding to  $\sim$ -11 eV, but from there to the Fermi level there are two strong antibonding orbitals in the molecular model.

The fact that the topmost filled bands are strongly antibonding within a Te<sub>5</sub> unit suggests that oxidation of the chain would lead to shorter Te-Te bands within that unit. With the Fermi level at  $\sim$ -7 eV oxidation is not excessively difficult and all three chains with a formal building unit of Te<sub>5</sub><sup>x-</sup> (-2 < x < 0) might be attainable with proper choice of synthetic conditions. All three known compounds were obtained by heating stoichiometric quantities of the elements involved;<sup>1-3</sup> compounds such as Rb<sub>2</sub>Te<sub>5</sub> and K<sub>2</sub>SnTe<sub>5</sub> with the lower electron count on the Te<sub>5</sub> unit might result from starting materials in the proper stoichiometric ratios.

The symmetry properties of the chains aid in understanding the location of the cations with respect to the central Te of a Te<sub>5</sub> unit. In the cisoid and transoid all-Te chains 2 and 3 the cations (Cs and Rb, respectively) are located on axes perpendicular to the Te<sub>5</sub> unit. In 2 the Cs-Te distances are 3.92 and 3.96 Å while in 3 the Rb-Te distance is 3.88 Å.

In the  $K_2SnTe_5$  structure the K<sup>+</sup> cations lie on an axis that runs through the central Te but in preservation of the symmetry of the tetragonal space group *I4cm* is at a 45° angle to the Te<sub>5</sub> plane. The Te-K distance is 4.24 Å. In fact, the cation is closer to the two peripheral Te's with distances of 3.52, 3.59, 3.71, and 3.80 Å.

The difference in the geometry of cation arrangement around the chains may be understood on the basis of the symmetry properties of the chains in the three structures. In the cisoid chain the plane of the Te<sub>5</sub> unit is not a mirror plane of symmetry of the entire chain. Consequently the HOMO and the orbital below it mix to form a pair of hybrid sp type orbitals (**21a** and **21b**) at the central Te. Thus there is a concentration of electron density along this axis and, not unexpectedly, the two cations for each Te<sub>5</sub> unit are located on this axis. In this chain there must be an



asymmetric top-bottom distribution of axial density at  $Te_c$ —and this might account for the 0.04-Å difference in  $Te_c$ -Cs distances. The longer  $Te_c$ -Cs distance is to Cs(2), which caps the "boat" formed by a  $Te_5$  unit and the four  $Te_p$ 's bonded to it. The distances between these last four  $Te_p$ 's and Cs(2) are all equal and are the shortest Te-Cs distance (3.85 Å), resulting in  $C_{4v}$  local symmetry and suggesting a chelating effect of these  $Te_p$ 's. Indeed, one of the lone-pair lobes on these atoms does point toward the inside of the "boat" and is certainly sufficiently diffuse to interact with the Cs cation.

The presence of an inversion center at  $Te_c$  in the transoid chain prevents s,  $p_z$  mixing; however, it does require equal  $Te_c-Rb$ distances above and below the  $Te_5$  plane. Similarly, there is no physical basis for mixing as in **21** since the  $Te_5$  units lie on a mirror plane. There is less of a concentration of electron density on  $Te_c$ , and the alternate location of the cation is preferred.

Another interesting feature of the structures is the short Te--Te interchain distance of 3.47 Å found in the transoid structure. Short distances such as this are found in many tellurium and selenium structures where steric factors do not prevent the close approach of these atoms. For instance, in  $\alpha$ -tellurium, the helical chains have Te-Te bond lengths of 2.83 Å, but each tellurium has four neighbors at distances of 3.494 Å.<sup>22</sup>

The geometric nature of the interaction between transoid chains is shown in 22 and 23. 22 indicates in a side view the short



contact. In 23 the same structure is shown in projection onto the plane of a Te<sub>5</sub> unit. One atom has been darkened and another shaded to aid in identifying the interaction. 23 in particular strongly suggests an interaction between  $p_z$ -type orbitals.

Calculations were carried out on a "dimer" composed of two molecules of 14. For the geometry as shown in 22 the dimer is not bound with respect to the two monomers by 0.47 eV. Shifting the molecules closer to 3.27 Å leads to a worse situation where the difference is  $\sim 1$  eV. A lateral shift of the chains to bring the two interacting atoms into an eclipsed orientation does not

<sup>(22)</sup> Cherin, P.; Unger, P. Inorg. Chem. 1967, 6, 1589-1591.

Table II. Extended Hückel Parameters

 orbital		$H_{ii}$ , eV	exponent	
 Te	5s	-20.78	2.51	
	5p	-11.04	2.16	
Sn	5s	-16.16	2.13	
	5p	-8.32	1.82	
Н	1 <b>s</b>	-13.60	1.30	

lead to any significant improvement. Interestingly, the reduced overlap populations for atoms that might be involved in this interaction show that indeed there is a small bonding interaction (+0.0039) but there is an antibonding one of similar magnitude (-0.0037) between the darkened Te and the central Te in the Te<sub>5</sub> unit containing the shaded Te. These two interactions are of similar magnitude in all the other geometries. As we move the two trimers toward each other, past the equilibrium separation in the polymer, the Te<sub>p</sub>-Te<sub>p</sub>' interaction grows more bonding but the above-mentioned Te<sub>p</sub>-Te<sub>c</sub>' interaction becomes more antibonding. The total energy for the approach of the two trimers is quite repulsive. It appears that our method does not portray properly the origins of what we feel on structural grounds must be a bonding interaction.

In summary, we have examined the electronic structure and bonding in a number of polymers containing the square-planar Te<sub>5</sub> repeating unit. These are seen to be characterized by weaker electron-rich three-center bonds within the unit, which accounts for the elongation of the intraunit Te–Te bond length by  $\sim 0.25$ Å over a normal bond, and the unit is shown to be an unstable species, compared with, e.g., ICl<sub>4</sub><sup>-</sup>. However, the formation of infinite one-dimensional chains due to covalent or dative bonding between units (Te<sub>5</sub><sup>2-</sup>) or with another modulating atom (SnTe<sub>5</sub><sup>2-</sup>) is energetically sufficient to stabilize the overall structure. In spite of the significant topological difference between the  $Cs_2Te_5$  (cisoid) and  $Rb_2Te_5$  (transoid) chains their electronic structures are virtually identical. These have been interpreted for all three compounds by employing model molecules, which provided a basis for understanding the band structure of the infinite chains. The HOMO in the model compounds is an antibonding orbital within the square-planar unit, leading to the suggestion that modification of the synthesis of these compounds could lead to an oxidized form of the chain with modified properties and greater inherent stability of the  $Te_5^{n-}$  units. The electronic basis for the short interchain Te-..Te distance of 3.47 Å, reminiscent of a similar contact in elemental tellurium, was investigated but apparently is beyond the capability of the extended Hückel method.

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#### Appendix

Extended Hückel parameters for all atoms used are presented in Table II. The  $H_{ii}$  values are from ref 23, and the exponents are from ref 24.

**Registry No.**  $Rb_2Te_5$ , 88188-94-7;  $Cs_2Te_5$ , 83332-22-3;  $K_2SnTe_5$ , 86205-22-3; Te, 13494-80-9.

(23) Hinze, J.; Jaffé, H. H. J. Am. Chem. Soc. 1962, 84, 540-548.
 (24) Clementi, E.; Roetti, C. At. Data. Nucl. Data Tables 1974, 14, 177-478.

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# Tetrahedral and Square-Planar One-Dimensional Chains: The Interplay of Crystal Field and Bandwidth in MS<sub>2</sub> Compounds

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One-dimensional  $ML_2$  chains with edge-sharing tetrahedral or square-planar coordination at the transition metal are the subject of this theoretical analysis. The band structures of these systems are examined in substantial detail. The local ligand field at the metal is a good starting point for an understanding of the band ordering in these extended structures, but inter unit cell interactions, both through the bridges and directly between the metals, are important for setting the energy order and dispersion of the bands and especially for determining unexpected electron counts for semiconducting behavior; one such case is d<sup>6</sup> for the tetrahedral chain. Pairing distortions, chain folding, and uniform contractions and elongations are also studied in some detail for these materials.

The two most frequency encountered geometrical arrangements of four ligands L around a metal atom are tetrahedral (1) or square planar (2); the complex stoichiometry in both cases in ML<sub>4</sub>. If



the ligand L has the ability to function as a bridge between two metal centers, it is in principle possible to generate from 1 and 2 the polymers 3 and 4, respectively. The stoichiometry of these



chains is  $ML_{4/2}$  or  $ML_2$ . The magnificent world of structural

Table I. Some Known ML<sub>2</sub> Structures

compd	chain	type	$\theta$ , deg	M-M, Å	d″	ref
KFeS <sub>2</sub>	FeS <sub>2</sub> <sup>-</sup>	3	106.0	2.70	d5	2a
$RbFeS_2$	FeS <sub>2</sub> <sup>-</sup>	3	105.0	2.71	d5	2b
CsFeS <sub>2</sub>	$FeS_2^-$	3	105.0ª	2.71ª	d <sup>5</sup>	2b, 3
Na <sub>3</sub> Fe <sub>2</sub> S <sub>4</sub>	FeS <sub>2</sub> <sup>-1.5-</sup>	3	106.0	2.75	d <sup>5.5</sup>	4
Na <sub>5</sub> Co <sub>2</sub> S <sub>5</sub>	CoS <sub>2</sub> <sup>1.5-</sup>	3	96.0	3.11	d <sup>6.5</sup>	5
$Na_2PdS_2$	$PdS_{2}^{-2-}$	4	82.5	3.54	d <sup>8</sup>	6
Na <sub>2</sub> PtS <sub>2</sub>	PtS <sub>2</sub> <sup>2~</sup>	4	82.5	3.55	d <sup>8</sup>	6
$K_2PtS_2$	$PtS_{2}^{2-}$	4	81.0	3.59	d <sup>8</sup>	7
$Rb_2PtS_2$	$PtS_{2}^{2-}$	4	79.1	3.64	d <sup>8</sup>	7
$\alpha$ -PdCl <sub>2</sub>	PdC1 <sub>2</sub>	4	87.0	3.34	d <sup>8</sup>	8

<sup>a</sup> Average value.

solid-state chemistry shows us that infinite systems such as 3 and 4 are not just the fruit of our imagination: they do occur in reality.<sup>1</sup>

<sup>(1)</sup> Bronger, W. Angew. Chem. 1981, 93, 12; Angew. Chem., Int. Ed. Engl. 1981, 20, 52.